

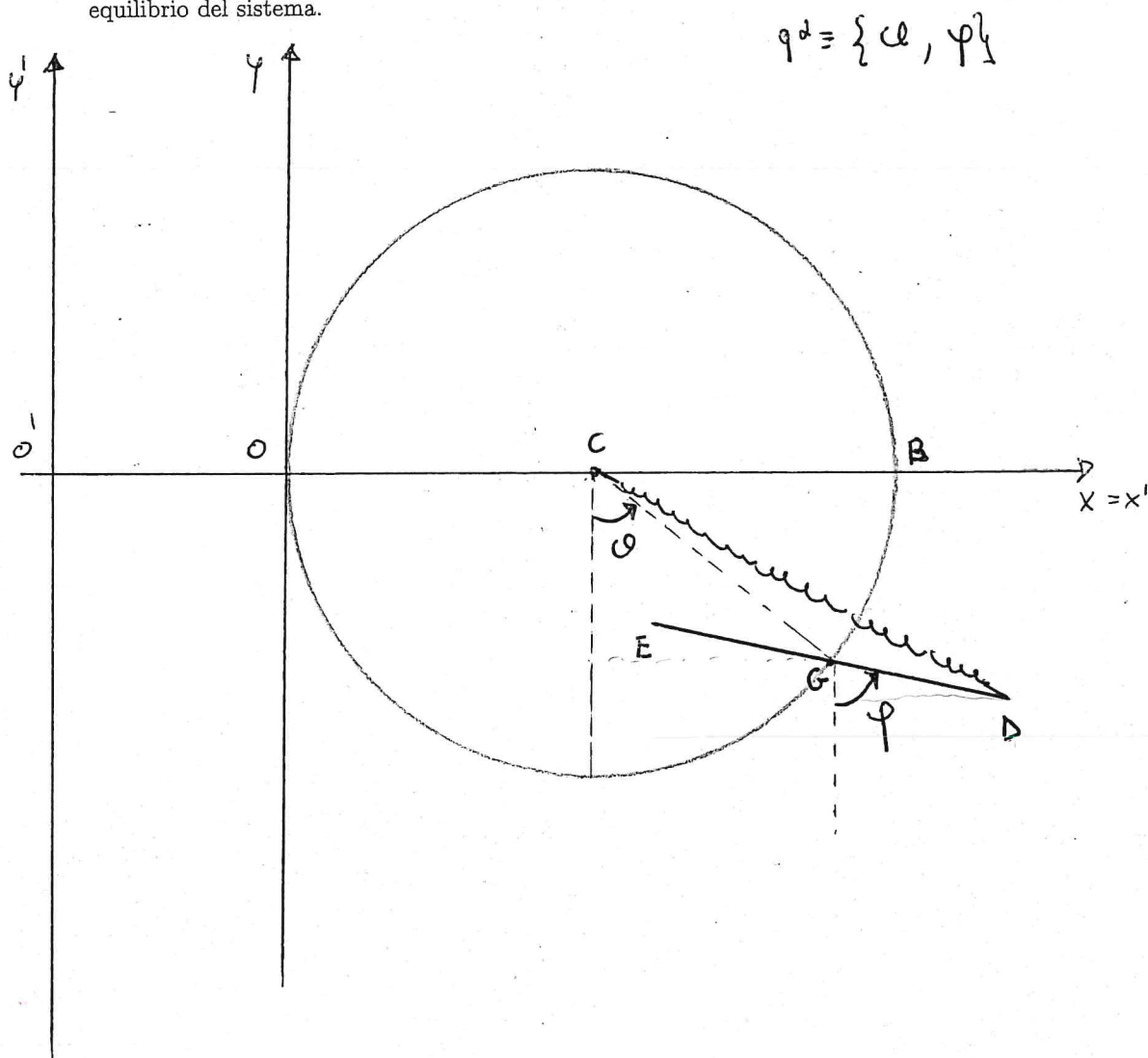
Università degli studi di Catania
 Corso di laurea Triennale in Fisica
 Prova scritta di Meccanica Analitica
 Appello del 02.12.2016

In un piano verticale, si consideri un riferimento fisso $\{o', x', y'\}$ ed un riferimento mobile relativo $\{o, x, y\}$, con gli assi orizzontali x ed x' costantemente sovrapposti, che si muove di moto traslatorio lungo la direzione x' con accelerazione costante $a = g$ (essendo g l'accelerazione di gravità). Solidale con il riferimento mobile relativo, sia dato un sistema S costituito da una circonferenza di raggio R , centro C e diametro OB sull'asse delle x (vedi figura), e da un'asta ED di lunghezza $2L$, massa m , avente densità non omogenea $\rho(s) = \alpha(L - s)^2$ (essendo s la distanza di un generico punto p dell'asta a partire da uno dei suoi estremi) con $0 \leq s \leq 2L$, ed α una costante positiva, il cui baricentro G può scorrere lungo la guida ~~in~~ circolare. Sul sistema oltre alla forza peso (lungo la verticale) agisce la forza elastica

$$\{F = -k(D - C), D\} \quad (1)$$

essendo k una costante positiva. Scegliendo le coordinate θ e ϕ come in figura, e supponendo tutti i vincoli lisci, si chiede di determinare nel riferimento mobile relativo:

1. Le configurazioni di equilibrio del sistema, discutendone la stabilità;
2. Le equazioni del moto e gli eventuali integrali primi del moto;
3. I moti in prima approssimazione attorno alle possibili configurazioni di equilibrio del sistema.



$$\text{Sia } \rho = a(L-x)^2 \quad 0 \leq x \leq 2L$$

(1)

$$m = \int_0^{2L} a(L-x)^2 dx = a \int_0^{2L} (L^2 + x^2 - 2Lx) dx =$$

$$= a \left\{ L^2 (2L) + \frac{(2L)^3}{3} - 2L \frac{(2L)^2}{2} \right\} =$$

$$= a \left\{ 2L^3 + \frac{8}{3} L^3 - \frac{8L^3}{2} \right\} = a L^3 \left\{ 2 + \frac{8}{3} - 4 \right\} =$$

$$= a L^3 \left\{ \frac{8}{3} - 2 \right\} = a \frac{2}{3} L^3 = m \Rightarrow$$

$$\boxed{a = \frac{3}{2} \frac{m}{L^3}}$$

CALCOLO DEL BARICENTRO



$$G = \frac{1}{m} \int_0^{2L} x \rho(x) dx =$$

$$= \frac{a}{m} \int_0^{2L} x (L-x)^2 dx =$$

$$= \frac{a}{m} \int_0^{2L} (L^2 x + x^3 - 2Lx^2) dx$$

$$= \frac{a}{m} \cdot \left\{ L^2 \frac{(2L)^2}{2} + \frac{(2L)^4}{4} - 2L \frac{(2L)^3}{3} \right\} =$$

$$= \frac{a}{m} \left\{ 2L^4 + 4L^4 - \frac{16L^4}{3} \right\} = \frac{a}{m} L^4 \left\{ 6 - \frac{16}{3} \right\} =$$

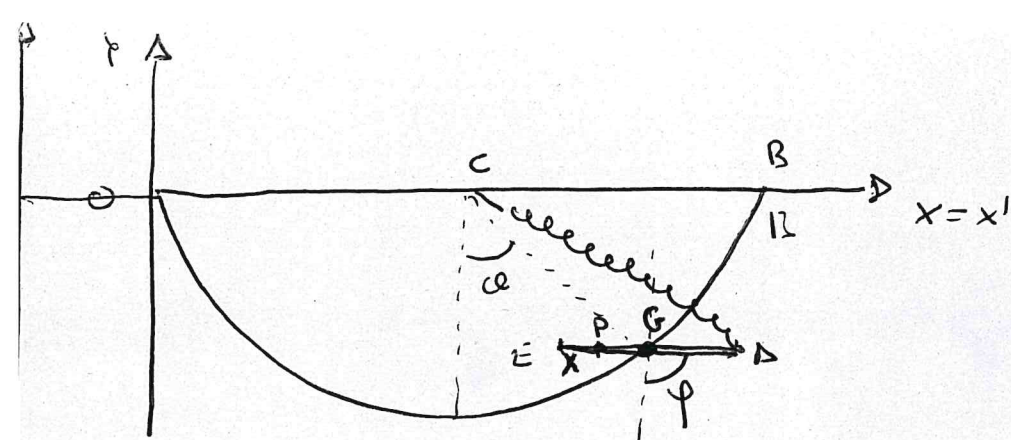
$$= \frac{a}{m} L^4 \cdot \frac{2}{3} = \frac{L^4}{m} \cdot \frac{2}{3} \cdot \left(\frac{3}{2} \frac{m}{L^3} \right) = L$$

ALORA IL BARICENTRO SI TROVA ANCORA NUC PUNTO MGAIO. (CORRE ASPETTARE, IN QUANTO $\rho(x)$ È SIMMETRICA RISPETTO AL PUNTO MGAIO) QUINDI:

$$E = \{ x_G - L \sin \varphi, -[|y_G| - L \cos \varphi] \} = \{ R \sin \alpha - L \sin \varphi, -R \cos \alpha + L \cos \varphi \}$$

$$D = \{ x_G + L \sin \varphi, -(|y_G| + L \cos \varphi) \} = \{ R \sin \alpha + L \sin \varphi, -R \cos \alpha - L \cos \varphi \}$$

(2)



$$\vec{G} \equiv \{ R \cos \varphi, -R \sin \varphi \} \quad C = [R, 0]$$

METODO: CALCOLO DELLA T UTILIZZANDO MONIE.

$$\bar{I}_{z,E} = \int_0^{2L} x^2 \rho(x) dx = a \int_0^{2L} x^2 (L-x)^2 dx =$$

$$= a \int_0^{2L} (L^2 x^2 + x^4 - 2L x^3) dx =$$

$$= a \left\{ L^2 \cdot \frac{(2L)^3}{3} + \frac{(2L)^5}{5} - 2L \frac{(2L)^4}{4} \right\} =$$

$$= a \left\{ \frac{8}{3} L^5 + \frac{32}{5} L^5 - \frac{32L^5}{4} \right\} = a L^5 \left\{ \frac{8}{3} + \frac{32}{5} - 8 \right\} =$$

$$= a L^5 \cdot \frac{40 + 96 - 120}{15} = \frac{16}{15} a L^5 = \frac{16}{15} \cdot \frac{3}{2} \frac{m}{L^2} L^5 = \frac{8}{5} m L^2$$

$$\bar{I}_{z,E} = \frac{8}{5} m L^2$$

$$\bar{I}_{z,E} = \bar{I}_{z,G} + m L^2 \Rightarrow \bar{I}_{z,G} = \bar{I}_{z,E} - m L^2$$

$$= \frac{8}{5} m L^2 - m L^2 =$$

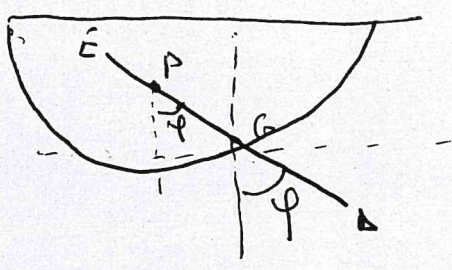
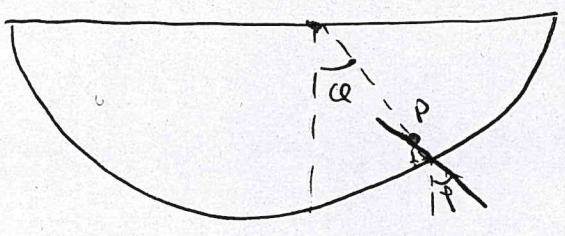
$$= \frac{3}{5} m L^2$$

$$T = \frac{1}{2} m \dot{\vec{G}}^2 + \frac{1}{2} \bar{I}_{z,G} \dot{\varphi}^2$$

$$\vec{G} = (R \cos \varphi, R \sin \varphi) \quad \dot{\vec{G}}^2 = R^2 \dot{\varphi}^2$$

$$T = \frac{1}{2} m R^2 \dot{\varphi}^2 + \frac{3}{10} m L^2 \dot{\varphi}^2$$

UTILIZZIAMO IL METODO DELLA INTEGRAZIONE PER IL CALCOLO DELLA T



$$\overline{PG} = \overline{EG} - x = (L - x)$$

$$P \equiv \left\{ x_G = (L-x) \sin \varphi, -[(y_G) = (L-x) \cos \varphi] \right\} =$$

$$\left\{ R \frac{1}{2} R \sin \alpha - (L-x) \sin \varphi, -R \cos \alpha + (L-x) \cos \varphi \right\}$$

$$\dot{P} = \left\{ R \cos \alpha \dot{\alpha} - (L-x) \cos \varphi \dot{\varphi}, R \sin \alpha \dot{\alpha} - (L-x) \sin \varphi \dot{\varphi} \right\}$$

$$\dot{P}^2 = R^2 \dot{\alpha}^2 + (L-x)^2 \dot{\varphi}^2 - 2R(L-x) \dot{\alpha} \dot{\varphi} [\cos \alpha \cos \varphi + \sin \alpha \sin \varphi]$$

DA cui

$$T = \frac{1}{2} R^2 \dot{\alpha}^2 \int dm + \frac{1}{2} \dot{\varphi}^2 \int_0^{2L} (L-x)^2 \rho(x) dx$$

$$= R [\cos \alpha \cos \varphi + \sin \alpha \sin \varphi] \dot{\alpha} \dot{\varphi} \int_0^{2L} (L-x) \rho(x) dx$$

$$= \frac{1}{2} m R^2 \dot{\alpha}^2 + \frac{1}{2} \dot{\varphi}^2 a \int_0^{2L} (L-x)^4 dx$$

$$- R [\cos \alpha \cos \varphi + \sin \alpha \sin \varphi] \dot{\alpha} \dot{\varphi} a \int_0^{2L} (L-x)^3 dx$$

poniamo $z = L-x$, $dz = -dx$

$$\int_0^{2L} (L-x)^4 dx = - \int^{-L} z^4 dz = \int_{-L}^L z^4 dz = \frac{2}{5} L^5$$

$$\int_0^{2L} (L-x)^3 dx = - \int_L^{-L} z^3 dz = \int_{-L}^L z^3 dz = 0$$

Quindi:

$$T = \frac{1}{2} m R^2 \dot{\theta}^2 + \frac{1}{2} \dot{\varphi}^2 \cdot \left(\frac{3}{2} \frac{m}{L^2} \right) \cdot \frac{2}{5} L^5 = 2$$

$$T = \frac{1}{2} m R^2 \dot{\theta}^2 + \frac{3}{10} m L^2 \dot{\varphi}^2$$

che e' lo stesso risultato ottenuto con König.

(CALCOLO DEL POTENZIALE)

OSSERVIAMO CHE SE CALCOLIAMO LE FORZE AGENTI

AVEREMO: FORZA PESO, FORZA ELASTICA, FORZA FITTIZIA DI TRASCINAMENTO

$$\{ \underline{F}_p = m \underline{g} (0, -1), G \}$$

$$\underline{a}_c = \ddot{\theta} = \underline{g} \cdot (1, 0)$$

$$\{ \underline{F}_e = -k (\Delta - c), \Delta \}$$

$$\underline{F}_c = -m \underline{a}_c = -m \underline{g} (1, 0)$$

$$\{ \underline{F}_c = m \underline{g} (-1, 0), G \}$$

$$U = \underline{U}_p + \underline{U}_e + \underline{U}_c = \int m \underline{g} (0, -1) \cdot (G - \underline{u}) - \frac{1}{2} k (\Delta - c)^2 + m \underline{g} (-1, 0) \cdot (G - \underline{u}) =$$

$$\Delta - c = \{ R \sin \alpha + L \sin \varphi, -R \cos \alpha - L \cos \varphi \}$$

$$\begin{aligned} (\Delta - c)^2 &= R^2 + L^2 - 2RL \sin \alpha \sin \varphi + 2RL \sin \alpha \sin \varphi \\ &+ 2RL \cos \alpha \cos \varphi = R^2 + L^2 + 2RL [\cos \alpha \cos \varphi + \sin \alpha \sin \varphi] \\ &= R^2 + L^2 + 2RL \cos (\alpha - \varphi) \end{aligned}$$

$$\Delta \text{ cui: } U = m \underline{g} R \cos \alpha - k R L \cos (\alpha - \varphi) - m \underline{g} R \sin \alpha + c$$

$$Q_\alpha = \frac{\partial U}{\partial \alpha} = -mgR \sin \alpha + KR L \sin(\alpha - \varphi) - mgR \cos \alpha$$

$$= -mgR (\sin \alpha + \cos \alpha) + KR L \sin(\alpha - \varphi)$$

$$Q_\varphi = \frac{\partial U}{\partial \varphi} = -KR L \sin(\alpha - \varphi)$$

$$Q_\alpha = KR L \sin(\alpha - \varphi) - mgR (\sin \alpha + \cos \alpha)$$

$$Q_\varphi = -KR L \sin(\alpha - \varphi)$$

CALCOLIAMO ADOTTANDO LE SOLLECITAZIONI ⇒ DA CUI LE POSIZIONI

NOTA: OSSERVIAMO CHE:

$$dF_z = -g(1, 0) dm \quad \Delta A \text{ cui}$$

$$dQ_\alpha^z = g(-1, 0) dm \cdot \frac{\partial P}{\partial \alpha}$$

$$dQ_\varphi^z = g(-1, 0) dm \cdot \frac{\partial P}{\partial \varphi}$$

$$P = [R + R \sin \alpha - (L-x) \sin \varphi, -R \cos \alpha + (L-x) \cos \varphi]$$

$$\frac{\partial P}{\partial \alpha} = [R \cos \alpha, R \sin \alpha] ; \quad \frac{\partial P}{\partial \varphi} = [-(L-x) \cos \varphi, -(L-x) \sin \varphi]$$

$$= -(L-x) [\cos \varphi, \sin \varphi]$$

$$dQ_\alpha^z = -gR \cos \alpha dm = -gR \cos \alpha f(x) dx$$

$$dQ_\varphi^z = (L-x) \cos \varphi dm = (L-x) \cos \varphi f(x) dx$$

DA CUI INTEGRANDO:

$$Q_\alpha^z = -mgR \cos \alpha$$

$$Q_\varphi^z = \cos \varphi \cdot a \int_0^L (L-x)^3 dx =$$

$$= \cos \varphi \cdot a \int_0^L z^3 dz = 0$$

Da cui il potenziale associato alle forze di trascinamento

$$U = \int Q_u^z du + c_1(\varphi) = -mgR \int \cos u du + c_1(\varphi) = -mgR \sin u + c_1(\varphi)$$

$$U = \int Q_\varphi^z d\varphi + c_2(u) \Rightarrow c_2(u) = -mgR \sin u$$

Avendo quindi ~~c_1(\varphi) = 0~~ $c_1(\varphi) = 0$ $U = -mgR \sin u$.

CALCOLO DEL POTENZIALE: PRIMA CHE LA FORZA DI TRASCINAMENTO PUO' ESSERE CALCOLATA SOTTO
ANZI CHE POTREMO CALCOLARE DIRETTAMENTE IL POTENZIALE

Alle forze di trascinamento

$$U_p^z = \int (-1, 0) dm \cdot (p-0) = -g [R \sin u - (L-x) \sin \varphi] dm$$

Da cui

$$U_p^z = -mgR \sin u + \int \sin \varphi \cdot \int_0^{2L} (L-x) g(x) dx = -mgR \sin u$$

Cio' come se la forza di trascinamento fosse diretta
applicata al centro, il che e' vero perche' e' una forza
costante infatti

$$U_p^z = \int g(-1, 0) dm \cdot (p-0) \Rightarrow U_z = g(-1, 0) \cdot \int (p-0) dm = m(g-0) = mg(-1, 0) \cdot (g-0)$$

come abbiamo fatto all'inizio.



A INIZIO CONI CHU:

$$\begin{cases} T = \frac{1}{2} m R^2 \dot{\varphi}^2 + \frac{3}{10} m L^2 \dot{\varphi}^2 \\ U = mgR(\cos u - \sin u) - KRL \cos(u - \varphi) \end{cases}$$

$$\begin{cases} Q_u = KRL \sin(u - \varphi) - mgR(\sin u + \cos u) \\ Q_\varphi = -KRL \sin(u - \varphi) \end{cases}$$

$$Q_{\varphi} = -KR \sin(\alpha - \varphi) = 0 \quad \alpha - \varphi = 0, \bar{u} \Rightarrow \varphi = \alpha \quad \boxed{\alpha = \varphi + \bar{u}}$$

$$Q_{\alpha} = 0 \Rightarrow \sin \alpha + \cos \alpha = 0 \Rightarrow \sin \alpha = -\cos \alpha \in$$

$$\Rightarrow \alpha = \frac{3}{4}\bar{u}, -\frac{\pi}{4}$$

Quindi: $\alpha = -\frac{\pi}{4} \quad \varphi = -\frac{\pi}{4} \quad e \quad \varphi = -\frac{\pi}{4} - \bar{u} = -\frac{5}{4}\bar{u} = \frac{3}{4}\bar{u}$

$$\alpha = \frac{3}{4}\bar{u} \quad \varphi = \frac{3}{4}\bar{u} \quad e \quad \varphi = \frac{3}{4}\bar{u} - \bar{u} = -\frac{\pi}{4}$$

$$\mathcal{R} \cdot (\alpha, \varphi) = \left\{ \begin{aligned} s_1 &\equiv \left(-\frac{\pi}{4}, -\frac{\pi}{4}\right), \quad s_2 \equiv \left(-\frac{\pi}{4}, \frac{3}{4}\bar{u}\right) \\ s_3 &\equiv \left(\frac{3}{4}\bar{u}, \frac{3}{4}\bar{u}\right), \quad s_4 \equiv \left(\frac{3}{4}\bar{u}, -\frac{\pi}{4}\right) \end{aligned} \right\}$$

$$\begin{cases} Q_{\varphi\varphi} = KR L \cos(\alpha - \varphi) \\ U_{\alpha\alpha} = KR L \cos(\alpha - \varphi) - m g R (\cos \alpha - \sin \alpha) \\ U_{\alpha\varphi} = U_{\varphi\alpha} = -KR L \cos(\alpha - \varphi) \end{cases}$$

Analisi: $U_{\varphi\varphi}|_{s_1, s_3} = KR L > 0$

$$\begin{cases} U_{\alpha\alpha}|_{s_3} = KR L + m g R \sqrt{2} \\ U_{\alpha\alpha}|_{s_1} = KR L - m g R \sqrt{2} \end{cases}$$

$$U_{\alpha\varphi}|_{s_1, s_2} = -KR L < 0$$

$$H_{s_1} = \begin{vmatrix} kRl - \sqrt{2}mgR & -kRl \\ -kRl & kRl \end{vmatrix} = -\sqrt{2}mgkLR^2 < 0$$

s_1 INSTABILE

$$H_{s_2} = \begin{vmatrix} kRl + \sqrt{2}mgR & -kRl \\ -kRl & kRl \end{vmatrix} = \sqrt{2}mgkLR^2 > 0$$

$U_{pp} > 0 \quad H > 0 \quad \text{NO MAX} \quad s_3 \text{ INSTABILE}$

$U_{pp}|_{s_2} = -kRl$

$U_{acc}|_{s_2} = -kRl + \sqrt{2}mgR$

$U_{acc}|_{s_3} = kRl$

$U_{pp}|_{s_2} < 0 \quad H > 0$
 $\text{MAX} \quad s_2 \text{ STABILE}$

$$H = \begin{vmatrix} -\sqrt{2}mgR - kRl & kRl \\ kRl & -kRl \end{vmatrix}$$

$$= (\cancel{kRl})^2 + \sqrt{2}mgkLR^2 - (\cancel{kRl})^2 = +\sqrt{2}mgkLR^2 > 0$$

$$\begin{cases} U_{acc}|_{s_4} = -kRl + \sqrt{2}mgR \\ U_{pp}|_{s_4} = -kRl \\ U_{acc}|_{s_4} = kRl \end{cases}$$

$$H = \begin{vmatrix} \sqrt{2}mgR - kRl & kRl \\ kRl & -kRl \end{vmatrix} = -\sqrt{2}mgkLR^2 < 0$$

$\text{NO MAX} \Rightarrow s_4 \text{ INSTABILE}$

$$T = \frac{1}{2} m R^2 \dot{\alpha}^2 + \frac{3}{10} m L^2 \dot{\varphi}^2$$

$$\frac{\Delta T}{\Delta \dot{\alpha}} = m R^2 \dot{\alpha}, \quad \frac{d}{dt} \frac{\Delta T}{\Delta \dot{\alpha}} = m R^2 \ddot{\alpha}, \quad \frac{\Delta T}{\Delta \dot{\alpha}} = 0$$

$$\frac{\Delta T}{\Delta \dot{\varphi}} = \frac{3}{5} m L^2 \dot{\varphi}, \quad \frac{d}{dt} \frac{\Delta T}{\Delta \dot{\varphi}} = \frac{3}{5} m L^2 \ddot{\varphi}, \quad \frac{\Delta T}{\Delta \dot{\varphi}} = 0$$

Ad cui:

$$1) \quad m R^2 \ddot{\alpha} = K R L \sin(\alpha - \varphi) - m g R (\sin \alpha + \cos \alpha) = Q_\alpha$$

$$2) \quad \frac{3}{5} m L^2 \ddot{\varphi} = -K R L \sin(\alpha - \varphi) = Q_\varphi$$

LE FORZE SONO CONSERVATIVE QUINDI COME INTEGRALI

PRIMO AVENDO SOLTANTO L'ENERGIA $E = T - U$

"MOTI IN PRIMA APPROSSIMAZIONE"

CONTINUA I MOTI APPROSSIMATI ATTORNO LA $S_2 \equiv \left(-\frac{\pi}{2}, \frac{3}{2}\pi \right)$

LINEARIZZAMO AURONO:

$$m R^2 \ddot{\alpha} = \cancel{Q_\alpha}|_{S_2} + \left. \frac{\Delta Q_\alpha}{\Delta \alpha} \right|_{S_2} \left(\alpha + \frac{\pi}{2} \right) + \left. \frac{\Delta Q_\alpha}{\Delta \varphi} \right|_{S_2} \left(\varphi - \frac{3}{2}\pi \right)$$

$$\frac{\Delta Q_\alpha}{\Delta \alpha} \Big|_{S_2} = \nabla_{\alpha\alpha} \Big|_{S_2} = -K R L - \sqrt{2} m g R$$

$$\frac{\Delta Q_\alpha}{\Delta \varphi} \Big|_{S_2} = \nabla_{\alpha\varphi} \Big|_{S_2} = K R L$$

AA Cui:

$$m R^2 \ddot{\varphi} = - (k R L + \sqrt{2} m g R) \left(\alpha + \frac{\pi}{4} \right) + k R L \left(\beta - \frac{3}{4} \pi \right)$$

AA 2^o condition

$$\Rightarrow \frac{3}{5} m L^2 \ddot{\varphi} = \cancel{U_{\varphi}} + \left. \frac{\partial Q_{\varphi}}{\partial \alpha} \right|_{S_2} \left(\alpha + \frac{\pi}{4} \right) + \left. \frac{\partial Q_{\varphi}}{\partial \beta} \right|_{S_2} \left(\beta - \frac{3}{4} \pi \right)$$

$$\left. \frac{\partial Q_{\varphi}}{\partial \alpha} \right|_{S_2} = U_{\varphi \alpha} \Big|_{S_2} = k R L$$

$$\left. \frac{\partial Q_{\varphi}}{\partial \beta} \right|_{S_2} = U_{\varphi \beta} \Big|_{S_2} = -k R L$$

$$\frac{3}{5} m L^2 \ddot{\varphi} = k R L \left(\alpha + \frac{\pi}{4} \right) - k R L \left(\beta - \frac{3}{4} \pi \right)$$

AA Cui

$$\begin{cases} m R^2 \ddot{\alpha} + (k R L + \sqrt{2} m g R) \left(\alpha + \frac{\pi}{4} \right) - k R L \left(\beta - \frac{3}{4} \pi \right) = 0 \\ \frac{3}{5} m L^2 \ddot{\varphi} - k R L \left(\alpha + \frac{\pi}{4} \right) + k R L \left(\beta - \frac{3}{4} \pi \right) = 0 \end{cases}$$

AA Cui

$$\alpha + \frac{\pi}{4} = \alpha_0 e^{\lambda t} \quad \beta - \frac{3}{4} \pi = \beta_0 e^{\lambda t}$$

$$\left[m R^2 \lambda^2 + (k R L + \sqrt{2} m g R) \right] \alpha_0 - k R L \beta_0 = 0$$

$$\left[\frac{3}{5} m L^2 \lambda^2 + k R L \right] \beta_0 - k R L \alpha_0 = 0$$

$$k R L \left[m R^2 \lambda^2 + (k R L + \sqrt{2} m g R) \right] - k R L \left[\frac{3}{5} m L^2 \lambda^2 + k R L \right] = 0$$

$$m k R L \left[R^2 + \frac{3}{5} L^2 \right]$$

$$\left(\frac{3}{5} m^2 R^2 L^2 \right) \lambda^4 + \left[m k R^3 L + \frac{3}{5} m L^2 (k R L + \sqrt{2} m g R) \right] \lambda^2 + (k R L)^2 - (k R L)^2 = 0$$

ΚΥΡΙΟΣ ΟΥΚΑΙ Λ'ΕΞΩΛΩΘΗΚΕ (ΡΩΜΕΛΛΟ $\lambda^2 = z$)

$$a z^2 + b z + c = 0$$

ΣΕ ΠΕΡΙΦΕΡΕΙΑ ΟΥΚΕ (ΥΠΕΡΒΟΛΗ)
 $\Delta = b^2 - 4ac > 0$ ΑΥΡΟΝΟ
ΡΑΔΙΕΙ ΡΕΑΛΙ, ΕΠΙΧΡΗΣΙΜΟΝ
ΛΑ ΡΕΦΟΡΑ ΜΙ ΚΑΡΤΕΣΙΟ

$$\text{C.S.U.} \begin{cases} a = \frac{3}{5} m^2 R^2 L^2 > 0 \\ b = m k R^2 L + \frac{3}{5} m L^2 (k R L + \sqrt{2} m g R) > 0 \\ c = \sqrt{2} m g k L R^2 > 0 \end{cases}$$

ΔΑ ΟΥΙ ΕΙΣΕΓΑΝ ΟΥΚ ΑΥΚ ΡΑΔΙΕΙ

$$z_1, z_2 \text{ ΤΑΛΙ ΟΥΚΕ } z_1 + z_2 = -\frac{b}{a} < 0 \quad z_1 \cdot z_2 = \frac{c}{a} > 0$$

$$z_1, z_2 < 0 \quad \Delta \text{ ΟΥΙ } \lambda_{1,2} = \pm i \sqrt{|z_1|} \quad \lambda_{3,4} = \pm i \sqrt{|z_2|}$$

ΟΥΚΑΙ ΟΜΠΙΟ ΣΙΩΝΟΚ ΑΙ ΡΟΤΙ ΑΚΡΟΝΟΙ.

ΟΜΠΙΟ ΣΙΩΝΟΚ Ι ΡΟΤΙ ΑΠΡΟ ΣΣΙΜΑΤΙ ΑΤΙΟΡΑΟ ΔΑ $S_1 = (-\frac{\pi}{4}, -\frac{\pi}{4})$

ΔΑ ΟΥΙ ΛΙΝΕΑΡΙΟΚΑΝΟ ΑΥΡΟΝΟ?

$$m R^2 \ddot{\varphi} = \cancel{Q_a}|_{S_1} + \frac{\partial Q_a}{\partial \alpha} \Big|_{S_1} (\alpha + \frac{\pi}{4}) + \frac{\partial Q_a}{\partial \varphi} \Big|_{S_1} (\varphi + \frac{\pi}{4})$$

$$\frac{\partial Q_a}{\partial \alpha} \Big|_{S_1} = k R L - m g R \sqrt{2}$$

$$\frac{\partial Q_a}{\partial \varphi} \Big|_{S_1} = -k R L$$

ΔΑ ΟΥΙ

$$m R^2 \ddot{\varphi} = (k R L - m g R \sqrt{2}) (\alpha + \frac{\pi}{4}) - k R L (\varphi + \frac{\pi}{4})$$

ΔΑ ΛΛΑ 2^ο ΟΜΠΙΟ ΣΙΩΝΟΚ

$$\frac{3}{5} m L^2 \ddot{\varphi} = \cancel{Q_a}|_{S_1} + \frac{\partial Q_a}{\partial \alpha} \Big|_{S_1} (\alpha + \frac{\pi}{4}) + \frac{\partial Q_a}{\partial \varphi} \Big|_{S_1} (\varphi + \frac{\pi}{4})$$

$$\frac{\partial Q_a}{\partial \alpha} \Big|_{S_1} = -k R L$$

$$\frac{\partial Q_a}{\partial \varphi} \Big|_{S_1} = k R L$$

SEMPLIFICHIAMO I COEFF. DEL TRINOMIO

ROTTO

$$\bar{a} z^2 + \bar{b} z + \bar{c} = 0$$

$$\begin{cases} \bar{a} = \frac{3}{5} m R L > 0 \\ \bar{b} = \kappa R^2 + \frac{3}{5} L (\kappa L + \sqrt{2} m g) > 0 \\ \bar{c} = \sqrt{2} g \kappa R \end{cases}$$

$$\bar{\Delta} = \bar{b}^2 - 4\bar{a}\bar{c} = \left[\kappa R^2 + \frac{3}{5} L (\kappa L + \sqrt{2} m g) \right]^2 - \frac{12}{5} \sqrt{2} m g \kappa L R^2$$

$$= (\kappa R^2)^2 + \underbrace{\left(\frac{3}{5} L \right)^2 (\kappa L + \sqrt{2} m g)^2 + \frac{6}{5} \kappa R^2 L (\kappa L - \sqrt{2} m g)}_{(\kappa L - \sqrt{2} m g)^2 + 4\sqrt{2} m g \kappa L}$$

$$= \underbrace{\left[\kappa R^2 + \frac{3}{5} L (\kappa L - \sqrt{2} m g) \right]^2}_{> 0} + \underbrace{\frac{36}{25} \sqrt{2} m g \kappa L^3}_{> 0} > 0$$

DA CUI IL $\bar{\Delta}$ DEL TRINOMIO $\bar{a} z^2 + \bar{b} z + \bar{c} = 0$ È TALMENTE

CHÉ $\bar{\Delta} > 0 \Rightarrow$ RADICI REALI.

DA cui:

$$\frac{3}{5} m L^2 \ddot{\varphi} = -k R L \left(\alpha + \frac{u}{a} \right) + k R L \left(\varphi + \frac{u}{a} \right)$$

DA cui LE DUE EQUAZ. LINEARI

$$m R^2 \ddot{\alpha} + (m g R \sqrt{2} - k R L) \left(\alpha + \frac{u}{a} \right) + k R L \left(\varphi + \frac{u}{a} \right) = 0$$

$$\frac{3}{5} m L^2 \ddot{\varphi} + k R L \left(\alpha + \frac{u}{a} \right) - k R L \left(\varphi + \frac{u}{a} \right) = 0$$

PERO' SOLUZIONI DEL TIPO $\alpha + \frac{u}{a} = \alpha_0 e^{\lambda t}$ $\varphi + \frac{u}{a} = \varphi_0 e^{\lambda t}$

DA cui IL SISTEMA:

$$\left[m R^2 \lambda^2 + (m g R \sqrt{2} - k R L) \right] \alpha_0 + (k R L) \varphi_0 = 0$$

$$(k R L) \alpha_0 + \left[\frac{3}{5} m L^2 \lambda^2 - k R L \right] \varphi_0 = 0$$

DA cui L'EQUAZIONE:

$$\left(\frac{3}{5} m^2 R^2 L^2 \right) \lambda^4 + \left[-m k R^3 L + \frac{3}{5} m L^2 (m g R \sqrt{2} - k R L) \right] \lambda^2 - k R L (m g R \sqrt{2} - k R L) - (k R L)^2 = 0$$

$$-k R L (m g R \sqrt{2} - k R L) - (k R L)^2 = 0$$

$$\left(\frac{3}{5} m^2 R^2 L^2 \right) \lambda^4 + m L \left[\frac{3}{5} L (m g R \sqrt{2} - k R L) - k R^3 \right] \lambda^2 - m g k L R^2 \sqrt{2} = 0$$

$-m g k L R^2 \sqrt{2} = 0$ \Leftrightarrow (in questo caso e' banale verificare che $\Delta > 0$)

DA cui L'EQUAZIONE $a z^2 + b z + c = 0$ con $a > 0$

$c < 0$

DA cui ESSENDO $z_1 - z_2 = \frac{c}{a} < 0$ AURAMO CERTAMENTE

UNA RADICE POSITIVA O ALTRA NEGATIVA DA CUI AD OGGI

$z_1 > 0$ e $z_2 < 0 \Rightarrow \lambda_{1,2} = \sqrt{z_1}$ MOTI IMPROPRIO

NOTI APPROSSIMATE ATTORNO AA $\alpha_2 = (\frac{3}{4}\bar{a}, \frac{7}{4}\bar{a})$

c)

$$m R^2 \ddot{\alpha} = \cancel{Q_{\alpha}}|_{\alpha_2} + \frac{\partial Q_{\alpha}}{\partial \alpha}|_{\alpha_2} (\alpha - \frac{3}{4}\bar{a}) + \frac{\partial Q_{\alpha}}{\partial \varphi}|_{\alpha_2} (\varphi - \frac{3}{4}\bar{a})$$

$$= (K R L + m g R \sqrt{2}) (\alpha - \frac{3}{4}\bar{a}) + K R L (\varphi - \frac{3}{4}\bar{a})$$

$$\frac{3}{5} m L^2 \ddot{\varphi} = \cancel{Q_{\varphi}}|_{\alpha_2} + \frac{\partial Q_{\varphi}}{\partial \alpha}|_{\alpha_2} (\alpha - \frac{3}{4}\bar{a}) + \frac{\partial Q_{\varphi}}{\partial \varphi}|_{\alpha_2} (\varphi - \frac{3}{4}\bar{a}) =$$

$$= -K R L (\alpha - \frac{3}{4}\bar{a}) + K R L (\varphi - \frac{3}{4}\bar{a})$$

DA CUI LE DUE EQUAZIONI:

$$\begin{cases} m R^2 \ddot{\alpha} - (K R L + m g R \sqrt{2}) (\alpha - \frac{3}{4}\bar{a}) + K R L (\varphi - \frac{3}{4}\bar{a}) = 0 \\ \frac{3}{5} m L^2 \ddot{\varphi} + (K R L) (\alpha - \frac{3}{4}\bar{a}) - K R L (\varphi - \frac{3}{4}\bar{a}) = 0 \end{cases}$$

CORCO SOLUZIONI $\alpha - \frac{3}{4}\bar{a} = \alpha_0 e^{\lambda t}$ $\varphi - \frac{3}{4}\bar{a} = \varphi_0 e^{\lambda t}$

DA CUI IL SISTEMA:

$$\left[m R^2 \lambda^2 - (K R L + m g R \sqrt{2}) \right] \alpha_0 + K R L \varphi_0 = 0$$

$$K R L (\alpha_0) + \left[\frac{3}{5} m L^2 \lambda^2 - K R L \right] \varphi_0 = 0$$

DA CUI L'EQUAZIONE AGLI AUTORE:

$$\left(\frac{3}{5} m^2 R^2 L^2 \right) \lambda^4 + \left[-m K R^3 L - \frac{3}{5} m L^2 (K R L + m g R \sqrt{2}) \right] \lambda^2$$

$$+ K R L (K R L + m g R \sqrt{2}) - (K R L)^2 = 0$$

DA CUI L'EQUAZIONE AGLI TIPO $a z^2 + b z + c = 0$

$$a = \frac{3}{5} m^2 R^2 L^2 > 0 \quad b = - \left\{ \frac{3}{5} m L^2 (K R L + m g R \sqrt{2}) + m K R^3 L \right\} < 0$$

NOTA: IN OGNI CASO
CHÉ $\Delta > 0$ SI PUVI
COME NEL CASO a)

ΔA cui:

$$z_1 + z_2 = -\frac{b}{a} > 0 \quad ; \quad z_1 \cdot z_2 = \frac{c}{a} > 0$$

AUROMO QUINDI: DUE SOLUZIONI $z_1 \neq 0$ $z_2 \neq 0$

QUINDI: NOTI I PARABOLICI

$$\lambda_{1,2} = \pm \sqrt{z_1} \quad \lambda_{3,4} = \pm \sqrt{z_2}$$

NOTI APPROSSIMATI ATTORNO AA $S_4 \equiv \left(\frac{3}{4}\bar{u}, -\frac{\bar{u}}{4} \right)$

$$mR^2 \ddot{\psi} = \cancel{Q_{\psi}}|_{S_4} + \frac{\partial Q_{\psi}}{\partial \psi}|_{S_4} \left(\psi - \frac{3}{4}\bar{u} \right) + \frac{\partial Q_{\psi}}{\partial \varphi}|_{S_4} \left(\varphi + \frac{\bar{u}}{4} \right)$$

$$= (-kRL + \sqrt{2} m g R) \left(\psi - \frac{3}{4}\bar{u} \right) + kRL \left(\varphi + \frac{\bar{u}}{4} \right)$$

$$\frac{3}{5} m L^2 \ddot{\varphi} = \cancel{Q_{\varphi}}|_{S_4} + \frac{\partial Q_{\varphi}}{\partial \varphi}|_{S_4} \left(\varphi - \frac{3}{4}\bar{u} \right) + \frac{\partial Q_{\varphi}}{\partial \psi}|_{S_4} \left(\psi + \frac{\bar{u}}{4} \right) =$$

$$= kRL \left(\varphi - \frac{3}{4}\bar{u} \right) - kRL \left(\psi + \frac{\bar{u}}{4} \right)$$

$$\begin{cases} mR^2 \ddot{\psi} + (-\sqrt{2} m g R + kRL) \left(\psi - \frac{3}{4}\bar{u} \right) - kRL \left(\varphi + \frac{\bar{u}}{4} \right) = 0 \\ \frac{3}{5} m L^2 \ddot{\varphi} - kRL \left(\varphi - \frac{3}{4}\bar{u} \right) + kRL \left(\psi + \frac{\bar{u}}{4} \right) = 0 \end{cases}$$

COME SOLUZIONI $\psi - \frac{3}{4}\bar{u} = \psi_0 e^{\lambda t}$ $\varphi + \frac{\bar{u}}{4} = \varphi_0 e^{\lambda t}$

ΔA cui:

$$\left[mR^2 \lambda^2 + (kRL - \sqrt{2} m g R) \right] \psi_0 - kRL \varphi_0 = 0$$

$$-kRL \psi_0 + \left[\frac{3}{5} m L^2 \lambda^2 + kRL \right] \varphi_0 = 0$$

$$\Rightarrow \left[\frac{3}{5} m^2 R^3 L^2 \right] \lambda^4 + \left[m k R^3 L + \frac{3}{5} m L^2 (kRL - \sqrt{2} m g R) \right] \lambda^2$$

$$+ \kappa R L (\kappa R L - \sqrt{2} m g R) - (\kappa R L)^2 = 0$$

DA cui

$$a z^2 + b z + c = 0 \quad \text{con}$$

$$a = \frac{3}{5} m^2 R^2 L^2 > 0$$

(il discriminante $\Delta > 0$ significa
due radici reali)

$$b = m L \left\{ \kappa R^3 + \frac{3}{5} L (\kappa R L - \sqrt{2} m g R) \right\}$$

$$c = -\sqrt{2} m g \kappa R^2 L < 0$$

DA cui:

$$z_1 + z_2 = -\frac{b}{a} \quad \text{ed} \quad z_1 \cdot z_2 = +\frac{c}{a} < 0$$

Quindi: avremo una soluzione positiva e una negativa.

La soluzione sia $z_1 > 0$ e $z_2 < 0$ DA cui

$$\lambda_1, \lambda_2 = \pm \sqrt{z_1} \quad \text{moti iperbolici}$$

$$\lambda_3, \lambda_4 = \pm i \sqrt{|z_2|} \quad \text{moti armonici}$$

