

Università degli studi di Catania  
 Corso di laurea in Fisica  
 Compito di Meccanica Analitica  
 Appello del 12.02.2016

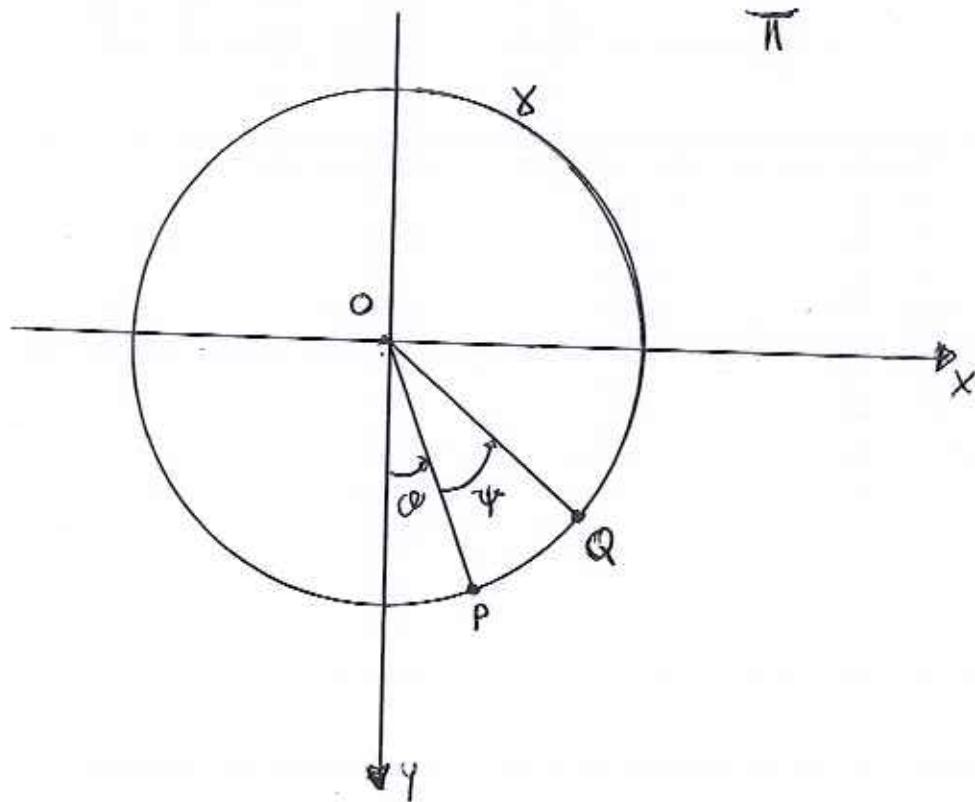
Sia data una guida circolare  $\gamma$  di centro  $O$  e raggio  $R$  posta in un piano verticale  $\Pi$  dove è stato introdotto un sistema di riferimento cartesiano ortogonale  $\{O, x, y\}$  con l'asse delle  $y$  verticale discendente. Un sistema materiale  $S$  è costituito da due punti  $P$  e  $Q$  aventi la stessa massa  $m$ , vincolati a muoversi senza attrito su  $\gamma$ , ed è soggetto, oltre alla forza peso, alle due forze di mutua repulsione

$$\{F, P\} \quad \text{e} \quad \{-F, Q\} \quad \text{con} \quad F = 4\alpha mgR \frac{(P - Q)}{|P - Q|^2}$$

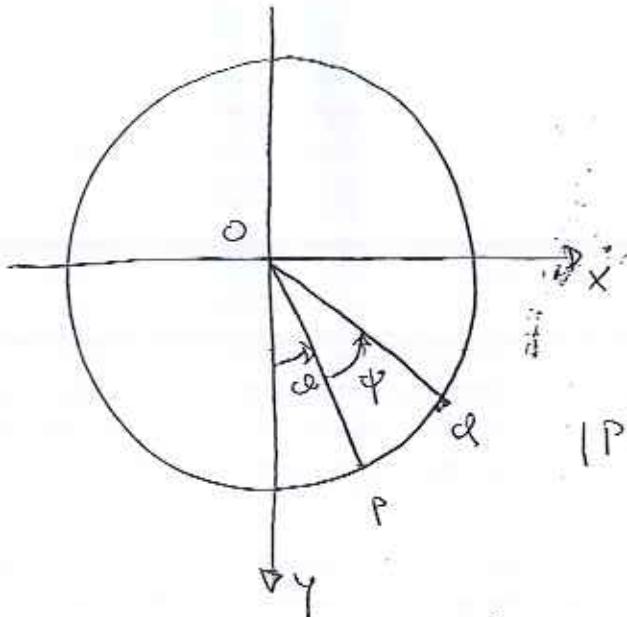
essendo  $\alpha$  una costante reale positiva.

Il sistema ha ovviamente due gradi di libertà, scelte allora come coordinate Lagrangiane gli angoli  $\vartheta$  che  $(P - O)$  forma con l'asse delle  $y$  e  $\psi$  che  $(Q - O)$  forma con  $(P - O)$  ambedue misurati in modo che le rotazioni di  $P$  per  $\vartheta$  crescente, e, per fissato  $P$ , quella di  $Q$  per  $\psi$  crescente siano entrambe in senso antiorario.

1. Dimostrare che la sollecitazione agente su  $S$  è conservativa e che per la coppia  $\{\vartheta, \psi\}$  delle due variabili lagrangiane  $0 \leq \vartheta \leq 2\pi$  e  $0 < \psi < 2\pi$ .
2. Determinare le configurazioni di equilibrio del sistema  $S$ , studiando la stabilità delle suddette configurazioni.
3. Determinare le equazioni di moto e gli eventuali integrali primi.
4. Studiare i moti linearizzati, determinando la frequenza dei piccoli moti, attorno ad una configurazione di equilibrio stabile.



①



$$P = [R \cos \alpha, R \sin \alpha]$$

$$Q = [R \cos(\alpha + \psi), R \sin(\alpha + \psi)]$$

$$|P-Q|^2 = R^2 [\cos \alpha - \cos(\alpha + \psi)]^2 +$$

$$+ R^2 [\sin \alpha - \sin(\alpha + \psi)]^2 =$$

$$= R^2 \{ \sin^2 \alpha + \sin^2(\alpha + \psi) - 2 \sin \alpha \sin(\alpha + \psi)$$

$$+ \cos^2 \alpha + \cos^2(\alpha + \psi) - 2 \cos \alpha \cos(\alpha + \psi) \}$$

$$= 2 R^2 \{ 1 - [\sin \alpha \sin(\alpha + \psi) + \cos \alpha \cos(\alpha + \psi)] \}$$

$\underbrace{\cos[(\alpha + \psi) - \alpha]}_{\cos \psi} = \cos \psi$

$$\Rightarrow |P-Q|^2 = 2 R^2 [1 - \cos \psi]$$

Proviamo che la somma rispetto su  $\psi$  è conservativa, proviamo

$$\text{cioè la forza } F = 4 \alpha mg R \frac{(P-Q)}{R^2} \quad \text{con } g = |P-Q|$$

è conio sui punti  $P, Q$   $\{F, P\} \subset \{-F, Q\}$

$\psi$ ' conservativa:

Ricordiamo che

$$F = f(g) \frac{(P-Q)}{g} \quad \text{mentre percorrendo } U = \iint f(g) d\gamma$$

$$\text{da cui essendo } f(g) = \frac{4 \alpha mg R}{g} \Rightarrow U = \iint \frac{4 \alpha mg R}{g} d\gamma =$$

$$= 4 \alpha mg R \ln(g) = 4 \alpha mg R \ln[V^2 R (1 - \cos \psi)^{1/2}] + C$$

$$= 2 \alpha mg R \ln(1 - \cos \psi) + K_1$$

Quindi:

$$U_F = 22mgR \ln(1 - \cos\psi) + U_L$$



NOTA: ALTRIMENTI SE CALCOLEREMO IL MURO

$$dL_{\text{muro}} = 42mgR \frac{(1-\alpha)}{g^2} \cdot d(1-\alpha) = 22mgR \frac{d\alpha^2}{g^2} =$$

$$= 22mgR \frac{d(1 - \cos\psi)}{(1 - \cos\psi)} = d[22mgR \ln(1 - \cos\psi)]$$

da cui  $\int dL = 0 \Rightarrow$  FORZA G' CONSERVATIVA



SOL. SISTEMA AURORA A NORD CO FOLGE POM

$$U_F + mg \cdot (p - \alpha) = mg(\alpha, z) \cdot (R \sin \omega, R \omega \sin \alpha) = mg R \cos \alpha$$

$$\begin{aligned} U_F &= mg \cdot (\alpha - \nu) = mg(\alpha, z) \cdot [R \sin(\alpha + \psi), R \cos(\alpha + \psi)] = \\ &= mg R \cos(\alpha + \psi) \end{aligned}$$

Quindi la potenziale totale è:

$$U_T = mgR \cos \alpha + mgR \cos(\alpha + \psi) + 22mgR \ln[1 - \cos\psi] + c$$

Proviamo a calcolare ciò

$$0 < \psi < 2\pi$$

INFATTI CONSIDERARE  $\dot{\theta} = \text{costante}$  (inizialmente zero)

$$\begin{aligned} T &= E + U \geq 0 \Rightarrow \begin{cases} \lim_{\psi \rightarrow 0} T = E + \lim_{\psi \rightarrow 0} U = -\infty \\ \lim_{\psi \rightarrow 2\pi} T = E + \lim_{\psi \rightarrow 2\pi} U = -\infty \end{cases} \text{ ALESSANDRO} \end{aligned}$$

Quindi avremo ciò

$$\begin{cases} 0 < \psi < 2\pi \\ 0 < \alpha < \pi \end{cases}$$

CALCOLARE CO SOLUZIONI DI ALTEZZE CONSERVATIVE

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$$Q_a = \frac{\sum \vec{F}}{\partial \alpha} = -mg R \sin \alpha - mg R \sin(\alpha + \psi)$$

$$Q_\psi = \frac{\sum \vec{F}}{\partial \psi} = -mg R \sin(\alpha + \psi) + 2\omega mg R \frac{\sin \psi}{1 - \cos \psi}$$

LE SUEZI RISULTANO VERSO CALCOLATO ANCHE DIRETTAMENTE.

ALLA CONSEGUENZA AVEMMO  $Q_a = \sum \underline{R_i}^{\text{for}} \cdot \frac{\Delta P_i}{\Delta \alpha}$ . INFATTI AVEMMO:

$$Q_a = mg \cdot \frac{\Delta \theta}{\Delta \alpha} + mg \cdot \frac{\Delta \phi}{\Delta \alpha} + f \cdot \frac{\Delta \phi}{\Delta \alpha} - F \frac{\Delta \phi}{\Delta \alpha} =$$

$$= mg(0,1) \cdot (R \cos \alpha, -R \sin \alpha) + mg(0,1) [R \cos(\alpha + \psi), -R \sin(\alpha + \psi)]$$

$$+ \frac{4\omega^2 mg R}{2R^2(1-\cos \psi)} \left\{ R [\sin \alpha - \sin(\alpha + \psi)], R [\cos \alpha - \cos(\alpha + \psi)] \right\} \cdot \{ R \cos \alpha, -R \sin \alpha\}$$

$$- \frac{4\omega^2 mg R}{2R^2(1-\cos \psi)} \cdot \left\{ R [\sin \alpha - \sin(\alpha + \psi)], R [\cos \alpha - \cos(\alpha + \psi)] \right\} \cdot \{ R \cos(\alpha + \psi), -R \sin(\alpha + \psi)\} =$$

$$= -mg R \sin \alpha - mg R \sin(\alpha + \psi) + \frac{4\omega^2 mg R}{2R^2(1-\cos \psi)} \left\{ R^2 \sin \alpha \cos \alpha - R^2 \cos \alpha \sin(\alpha + \psi) - R^2 \sin \alpha \sin(\alpha + \psi) + R^2 \cos \alpha \cos(\alpha + \psi) \right\}$$

$$- \frac{4\omega^2 mg R}{2R^2(1-\cos \psi)} \left\{ R^2 \sin \alpha \cos(\alpha + \psi) - R \sin \alpha \cos(\alpha + \psi) \sin(\alpha + \psi) - R^2 \cos \alpha \sin(\alpha + \psi) + R^2 \sin \alpha \cos(\alpha + \psi) \sin(\alpha + \psi) \right\}$$

$$= -mg R \sin \alpha - mg R \sin(\alpha + \psi) + \frac{2\omega^2 mg R}{(1-\cos \psi)} [\sin \alpha \cos(\alpha + \psi) - \cos \alpha \sin(\alpha + \psi)] - \frac{2\omega^2 mg R}{(1-\cos \psi)} [\sin \alpha \cos(\alpha + \psi) - \cos \alpha \sin(\alpha + \psi)]$$

$$\Rightarrow Q_a = -mg R \sin \alpha - mg R \sin(\alpha + \psi)$$

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$$\ddot{\varphi}_4 = mg \cdot \cancel{\frac{d\varphi}{d\psi}} + mg \cdot \frac{d\varphi}{d\psi} + F \cdot \cancel{\frac{d\varphi}{d\psi}} - F \cdot \frac{d\varphi}{d\psi}$$

$$= mg(0,1) \cdot [R \omega_2 (\alpha + \dot{\psi}), -R \sin(\alpha + \dot{\psi})]$$

$$- \frac{R \cdot mg R}{2R^2 [1 - \cos \dot{\psi}]} \cdot [R(\sin \alpha - \sin(\alpha + \dot{\psi})), R(\cos \alpha - \cos(\alpha + \dot{\psi}))]$$

$$\cdot [R \omega_2 (\alpha + \dot{\psi}), -R \sin \alpha (\alpha + \dot{\psi})]$$

$$= -mg R \sin(\alpha + \dot{\psi}) - \frac{4 \cdot mg R}{2R^2 (1 - \cos \dot{\psi})} \left\{ R^2 \sin \alpha \cos(\alpha + \dot{\psi}) \right.$$

$$- R^2 \sin(\alpha + \dot{\psi}) \cos(\alpha + \dot{\psi}) - R^2 \cos \alpha \sin(\alpha + \dot{\psi}) + R^2 \sin(\alpha + \dot{\psi}) \cos(\alpha + \dot{\psi})$$

$$= -mg R \sin(\alpha + \dot{\psi}) - \frac{2 \cdot mg R}{(1 - \cos \dot{\psi})} \left[ \text{since } \cos(\alpha + \dot{\psi}) - \cos \alpha \sin(\alpha + \dot{\psi}) \right]$$

$$\sin(\alpha - (\alpha + \dot{\psi})) =$$

$$= -2 \sin \dot{\psi}$$

$$\ddot{\varphi}_4 = -mg R \sin(\alpha + \dot{\psi}) + 2 \cdot mg R \frac{\sin \dot{\psi}}{1 - \cos \dot{\psi}}$$

"Educiendo"

$$\begin{aligned} \alpha &= 0 \Rightarrow \begin{cases} \text{since } +2 \sin(\alpha + \dot{\psi}) = 0 & (1) \\ -2 \sin(\alpha + \dot{\psi}) + 2 \cdot \frac{\sin \dot{\psi}}{1 - \cos \dot{\psi}} = 0 & (2) \end{cases} & \text{basissinarii} \\ \alpha &= \bar{\alpha} \Rightarrow \begin{cases} \alpha + \dot{\psi} = 2k\pi - \alpha \\ \alpha + \dot{\psi} = \bar{\alpha} - (-\alpha) \end{cases} & \text{b. caso } n=0, 1 \end{aligned}$$

$$\text{DAUCA (1)} \quad \sin(\alpha + \dot{\psi}) = -\sin(\alpha) = \sin(-\alpha) \Rightarrow \begin{cases} \alpha + \dot{\psi} = 2k\pi - \alpha \\ \alpha + \dot{\psi} = \bar{\alpha} - (-\alpha) \end{cases}$$

$$\text{DAUCA PRIMA} \Rightarrow \dot{\psi} = 2\pi - 2\alpha$$

OSSERVATO I VALORI  $\alpha = 0, \bar{\alpha}, 2\bar{\alpha}$  NON SONO ACCETTABILI se  $0 < \dot{\psi} < 2\bar{\alpha}$

$$\text{INFATTI} \begin{cases} \text{se } \alpha = 0 \quad \dot{\psi} = 2\bar{\alpha} \\ \text{se } \alpha = \bar{\alpha} \quad \dot{\psi} = 0 \\ \text{se } \alpha = 2\bar{\alpha} \quad \dot{\psi} = -2\bar{\alpha} \end{cases}$$

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Solutions occur if the condition  $\psi = 2\bar{u} - \alpha$ 

$$-\sin(2\bar{u} - \alpha) + 2d \frac{\sin(2\bar{u} - 2\alpha)}{1 - \cos(2\bar{u} - 2\alpha)} = 0$$

$$\sin(2\bar{u} - \alpha) = \sin(-\alpha) = -\sin\alpha$$

$$\sin(2\bar{u} - 2\alpha) = \sin(-2\alpha) = -\sin(2\alpha) = -2\sin\alpha\cos\alpha$$

$$\cos(2\bar{u} - 2\alpha) = \cos(-2\alpha) = \cos(2\alpha) = \cos^2\alpha - \sin^2\alpha$$

$$+\sin\alpha \cdot \frac{4d \text{ Since Case}}{1 - \cos^2\alpha + \sin^2\alpha} = \text{Since} - \frac{4d \text{ Hard Case}}{2\sin^2\alpha} =$$

$$= \text{Since} - 2d \frac{\cos\alpha}{\text{Since}} = 0 \Rightarrow \frac{\text{Since} - 2d \cos\alpha}{\text{Since}} = 0$$

$\Rightarrow$  (since since  $\neq 0$ )

$$\Rightarrow \cos^2\alpha + 2d \cos\alpha - 1 = 0 \quad \text{Position } \cos\alpha = x$$

$$\Rightarrow x^2 + 2d x - 1 = 0 \Rightarrow x = \cos\alpha = -d \pm \sqrt{d^2 + 1}$$

Considering the solution  $-d - \sqrt{d^2 + 1} < -1 \quad (-1 \leq \cos\alpha \leq 1)$

$$\text{Hence } \cos\alpha = -d + \sqrt{d^2 + 1} > 0 \quad \text{Hence } \alpha = \arccos(-d + \sqrt{d^2 + 1})$$

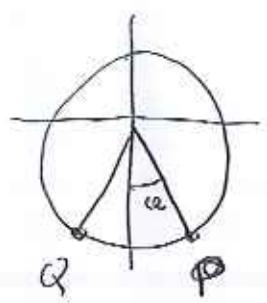
$$\text{Quint. } \Theta_1 = \arccos(-d + \sqrt{d^2 + 1}) = \bar{\alpha} \quad \psi_1 = 2\bar{u} - 2\bar{\alpha}$$

$$\Theta_2 = -\bar{\alpha} \quad \text{e} \quad \psi_2 = 2\bar{u} + 2\bar{\alpha} = \bar{\psi}$$

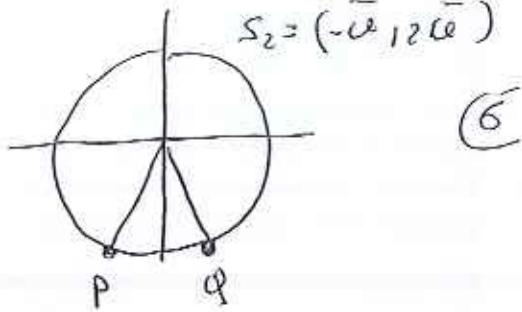
Position  $0 < \psi < 2\bar{u}$

$$S_1 = (\omega, \psi_1) = (\bar{\alpha}, 2\bar{u} - 2\bar{\alpha})$$

$$S_2 = (-\bar{\Theta}_2, \bar{\psi}_2) = (-\bar{\alpha}, 2\bar{\alpha})$$



$$S_1 = (\alpha, \bar{u} - 2\bar{\omega})$$



(6)

B)

Si considerano le altre 3 configurationi

$$\alpha + \psi = \bar{u} + \bar{\phi} \Rightarrow \bar{\psi} = \bar{u}$$

$$\text{DAGLA (2)} \quad \sin(\alpha + \bar{u}) = 0 \Rightarrow -\sin(\alpha) = 0 \Rightarrow \alpha = 0, \pi$$

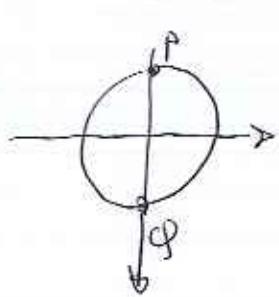
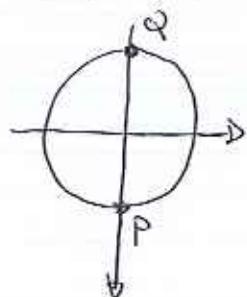
Altre 3 delle 6 configurazioni:

$$S_2 = (\alpha_2, \psi_2) = (0, \bar{u})$$

$$S_4 = (\alpha_4, \psi_4) = (\bar{u}, \bar{u})$$

$$S_1 = (0, \bar{u})$$

$$S_3 = ($$



Stabilimento della stabilità:

$$\frac{d\alpha}{d\alpha} = \frac{d^2 u}{d\alpha^2} = -mgR \cos \alpha - mgR \cos(\alpha + \psi)$$

$$\begin{aligned} \frac{d\psi}{d\psi} &= \frac{d^2 u}{d\psi^2} = -mgR \cos(\alpha + \psi) + 2\alpha mgR \underbrace{\frac{\cos \psi(1 - \cos \psi) - \sin \psi \cdot 2\alpha \sin \psi}{(1 - \cos \psi)^2}}_{\frac{(1 - \cos \psi)}{(1 - \cos \psi)^2}} \\ &= -mgR \cos(\alpha + \psi) - 2\alpha mgR \end{aligned}$$

$$= -mgR \left\{ \cos(\alpha + \psi) + \frac{2\alpha}{1 - \cos \psi} \right\}$$

$$\frac{d^2 u}{d\cos \psi} = -mgR \cos(\alpha + \psi)$$

Si calcolano questi quantità per i punti

A)  $\frac{\sum \vec{U}}{\sum \vec{\alpha}^2} \Big|_{S_1, S_2} = -mgR \cos(\bar{\alpha}) - mgR (\cos \bar{\alpha}) = -2mgR \cos \bar{\alpha} < 0$

Per cui:  $0 < \cos \bar{\alpha} \leq 1$

$$\frac{\sum \vec{U}}{\sum \vec{\psi}^2} \Big|_{S_1, S_2} = -mgR \left[ \cos \bar{\alpha} + \frac{2\lambda}{1 - \cos(2\bar{\alpha})} \right] < 0$$

$$\frac{\sum \vec{U}}{\sum \vec{\alpha} \vec{\psi}} \Big|_{S_1, S_2} = -mgR \cos(\bar{\alpha}) < 0$$

da cui:

$$H \Big|_{S_1, S_2} = 2(mgR)^2 \left[ \cos \bar{\alpha}^2 + \frac{2\lambda \cos(\bar{\alpha})}{1 - \cos(2\bar{\alpha})} \right] \neq (mgR)^2 \cos' \bar{\alpha} = \\ = (mgR)^2 \left[ \cos' \bar{\alpha} + \frac{4\lambda \cos(\bar{\alpha})}{1 - \cos(2\bar{\alpha})} \right] > 0$$

Quindi:  $S_1, S_2$  sono stabili massimi per  $V \Rightarrow S_1, S_2$  sono stabili

B)  $\frac{\sum \vec{U}}{\sum \vec{\alpha}^2} \Big|_{S_3, S_4} = -mgR (\pm 1) - mgR (\mp 1) = 0$

$$\frac{\sum \vec{U}}{\sum \vec{\psi}^2} \Big|_{S_3, S_4} = -mgR \left\{ \dots \right\}$$

$$\frac{\sum \vec{U}}{\sum \vec{\alpha} \vec{\psi}} \Big|_{S_3, S_4} = -mgR (\mp 1)$$

Quindi:  $H_{S_3, S_4} = -(mgR)^2 < 0$

Quindi non sono massimi  $\Rightarrow S_3, S_4$  sono instabili

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CALCULAZIONE DELL'ENERGIA CINETICA

$$\mathcal{T} = \frac{1}{2} m \dot{\vec{r}}^2 + \frac{1}{2} m \dot{\vec{Q}}^2$$

$$\dot{\vec{r}} = [R \dot{\phi} \cos(\alpha + \psi), -R \dot{\phi} \sin(\alpha + \psi)] \Rightarrow \dot{\vec{r}}^2 = R^2 \dot{\phi}^2$$

$$\dot{\vec{Q}} = [R (\ddot{\phi} + \dot{\psi}) \cos(\alpha + \psi), -R (\ddot{\phi} + \dot{\psi}) \sin(\alpha + \psi)]$$

$$\dot{\vec{Q}}^2 = R^2 (\dot{\phi} + \dot{\psi})^2 = R^2 (\dot{\phi}^2 + \dot{\psi}^2 + 2\dot{\phi}\dot{\psi})$$

$$\mathcal{T} = \frac{1}{2} m R^2 (2\dot{\phi}^2 + \dot{\psi}^2 + 2\dot{\phi}\dot{\psi}) = m R^2 [\dot{\phi}^2 + \dot{\phi}\dot{\psi} + \frac{1}{2} \dot{\psi}^2]$$

$$\frac{d\mathcal{T}}{dt} = m R^2 [2\dot{\phi} + \dot{\psi}] \Rightarrow \frac{d}{dt} \frac{d\mathcal{T}}{d\dot{\phi}} = m R^2 [2\ddot{\phi} + \ddot{\psi}]$$

$$\frac{d\mathcal{T}}{d\dot{\phi}} = 0 \quad \frac{d\mathcal{T}}{d\dot{\psi}} = 0$$

$$\frac{d\mathcal{T}}{d\dot{\psi}} = m R^2 [\ddot{\psi} + \ddot{\phi}] \Rightarrow \frac{d}{dt} \frac{d\mathcal{T}}{d\dot{\psi}} = m R^2 [\ddot{\psi} + \ddot{\phi}]$$

Due equazioni:

$$\left\{ \begin{array}{l} m R^2 [2\ddot{\phi} + \ddot{\psi}] = -m g R \{ \sin \alpha + \sin(\alpha + \psi) \} \\ m R^2 [\ddot{\phi} + \ddot{\psi}] = -m g R \{ \sin(\alpha + \psi) - 2 \frac{\sin \psi}{1 - \cos \psi} \} \end{array} \right.$$

- ✓

INIZIATURA PIANETI:

1) FORZA DI GRAZIAZIONE  $\vec{F} = \vec{T} - \vec{U} = m \vec{a}_{\text{forz}}$ 

2) NUOVA RELAZIONE DELL'ANALISI CICLICA:

3) LE EQUAZIONI SONO ACCOPPIATE.

- ✓

# (9)

Problema 10: Spina Approssimazione

Supponiamo che la confezione 2'000 g di equivalente statico  
 $S_1 = (\ddot{\alpha}, 2\ddot{\alpha} - 2\ddot{\psi}) = (\ddot{\alpha}, -2\ddot{\alpha})$  con  $\omega \ddot{\alpha} = -d + \sqrt{d^2+1}$

La linearizzazione del suo momento attivo.

Momenti, angoli, velocità  
 $S_1 = (\ddot{\alpha}, -2\ddot{\alpha}, 0, 0, 0, 0)$

$$2R\ddot{\alpha} + R\ddot{\psi} + g [ \sin \alpha + \sin(\alpha + \psi) ] = F(\alpha, \psi, \ddot{\alpha}, \ddot{\psi}, \ddot{\ddot{\alpha}}, \ddot{\ddot{\psi}}) = 0$$

$$R\ddot{\ddot{\alpha}} + R\ddot{\ddot{\psi}} + g \left[ \sin(\alpha + \psi) - 2d \frac{\sin \psi}{1 - \cos \psi} \right] = G(\alpha, \psi, \ddot{\alpha}, \ddot{\psi}, \ddot{\ddot{\alpha}}, \ddot{\ddot{\psi}}) = 0$$

$$F|_{S_1} = 0 \quad \frac{\partial F}{\partial \ddot{\alpha}} = 2R \quad \frac{\partial F}{\partial \ddot{\psi}} = R \quad \frac{\partial F}{\partial \ddot{\ddot{\alpha}}} = \frac{\partial F}{\partial \ddot{\ddot{\psi}}} = 0$$

$$\begin{aligned} \frac{\partial F}{\partial \alpha} \Big|_{S_1} &= g [\cos \alpha + \cos(\alpha + \psi)] \Big|_{S_1} = g [\cos(\bar{\alpha}) + \cos(-\bar{\alpha})] = 2g \cos \bar{\alpha} \\ &= 2g [\sqrt{d^2+1} - d] \end{aligned}$$

$$\frac{\partial F}{\partial \psi} = g \cos(\alpha + \psi) \Big|_{S_1} = g \cos(-\bar{\alpha}) = g \cos \bar{\alpha} = g [\sqrt{d^2+1} - d]$$

Dunque

$$2R\ddot{\alpha} + R\ddot{\psi} + 2g \underbrace{[\sqrt{d^2+1} - d]}_{K} (\alpha - \bar{\alpha}) + g \underbrace{[\sqrt{d^2+1} - d]}_{K} (\psi + 2\bar{\alpha}) = 0$$

$$G|_{S_1} = + \frac{1}{\sin(\bar{\alpha})} [\cos^2 \bar{\alpha} + 2d \cos(\bar{\alpha}) + 1] = 0$$

$$\frac{\partial G}{\partial \ddot{\alpha}} = R ; \quad \frac{\partial G}{\partial \ddot{\psi}} = R ; \quad \frac{\partial G}{\partial \ddot{\ddot{\alpha}}} = \frac{\partial G}{\partial \ddot{\ddot{\psi}}} = 0$$

$$\frac{\partial G}{\partial \alpha} = g \left\{ \cos(\alpha + \psi) \right\}_{\alpha=0} = g \cos(-\bar{\alpha}) = g \cos(\bar{\alpha})$$

$$\frac{\partial G}{\partial \psi} = g \left\{ \cos(\alpha + \psi) - 2d \frac{\cos \psi (1 - \cos \psi) - \sin \psi \sin \psi}{(1 - \cos \psi)^2} \right\}_{\alpha=0} =$$

$$\begin{aligned}
 \frac{\Delta G}{\Delta \psi} &= g \left\{ \cos(\alpha + \bar{\alpha}) + \frac{2\omega}{1 - \omega \bar{\alpha}} \right\}_{S_1} = g \left[ \cos(-\bar{\alpha}) + \frac{2\omega}{1 - \omega(2\bar{\alpha})} \right] \quad (10) \\
 &= g \left[ \cos(\bar{\alpha}) + \frac{2\omega}{1 - \omega^2(\bar{\alpha}) + 4m^2(\bar{\alpha})} \right] = g \left[ \cos(\bar{\alpha}) + \frac{2\omega}{2m^2(\bar{\alpha})} \right] = \\
 &= g \left[ \cos(\bar{\alpha}) + \frac{1}{1 - \omega^2(\bar{\alpha})} \right]
 \end{aligned}$$

Nota: escono  $\omega \bar{\alpha} = -\omega + \sqrt{\omega^2 H}$   $\Rightarrow \omega^2 \bar{\alpha} = \omega^2 + \omega^2 + 1 - 2\omega \sqrt{\omega^2 H} = 2\omega^2 + 1 - 2\omega \sqrt{\omega^2 H} \Rightarrow 1 - \omega^2 \bar{\alpha} = 2\omega [-\omega + \sqrt{\omega^2 + 1}] = 2\omega \cos(\bar{\alpha})$

Quindi:

$$\frac{\Delta G}{\Delta \psi} \Big|_{S_1} = g \left[ \cos(\bar{\alpha}) + \frac{1}{2\omega \cos(\bar{\alpha})} \right] = g \left[ \cos(\bar{\alpha}) + \frac{1}{2\omega \cos(\bar{\alpha})} \right]$$

da cui arrivò:

$$2R \ddot{\phi} + R \ddot{\psi} + [ \dot{\psi} g \cos \bar{\alpha} ] (\alpha - \bar{\alpha}) + (g \cos \bar{\alpha}) (\psi + 2\bar{\alpha}) = 0$$

$$R \ddot{\phi} + R \ddot{\psi} + (g \cos \bar{\alpha}) (\alpha - \bar{\alpha}) + g \left[ \cos \bar{\alpha} + \frac{1}{2\omega \cos \bar{\alpha}} \right] (\psi + 2\bar{\alpha}) = 0$$

da cui si ricavano le soluzioni

$$\psi + 2\bar{\alpha} = \psi_0 e^{\lambda t} ; \quad \alpha - \bar{\alpha} = \alpha_0 e^{\lambda t}$$

$$2R\lambda^2 \alpha_0 + R\lambda^2 \psi_0 + (2g \cos \bar{\alpha}) \alpha_0 + (g \cos \bar{\alpha}) \psi_0 = 0$$

$$R\lambda^2 \alpha_0 + R\lambda^2 \psi_0 + (g \cos \bar{\alpha}) \alpha_0 + g \left( \cos \bar{\alpha} + \frac{1}{2\omega \cos \bar{\alpha}} \right) \psi_0 = 0$$

da cui le due equazioni di cui arrivò nelle variabili  $\alpha_0, \psi_0$

$$2[R\lambda^2 + g \cos \bar{\alpha}] \alpha_0 + [R\lambda^2 + g \cos \bar{\alpha}] \psi_0 = 0$$

$$[R\lambda^2 + g \cos \bar{\alpha}] \alpha_0 + \left\{ R\lambda^2 + g \left[ \cos \bar{\alpha} + \frac{1}{2\omega \cos \bar{\alpha}} \right] \right\} \psi_0 = 0$$

(11)

Dati i condizioni scalari

$$2[R\lambda^2 + g \omega_1 \bar{u}] \left\{ R\lambda^2 + g \left[ \omega_1 \bar{u} + \frac{1}{2\omega_1 \bar{u}} \right] \right\} -$$

$$- [R\lambda^2 + g \omega_1 \bar{u}]^2 = 0$$

$$\Rightarrow [R\lambda^2 + g \omega_1 \bar{u}] \left\{ 2R\lambda^2 + 2g \left[ \omega_1 \bar{u} + \frac{1}{2\omega_1 \bar{u}} \right] - [R\lambda^2 + g \omega_1 \bar{u}] \right\} = 0$$

$$\Rightarrow \boxed{\lambda^2 = - \frac{g}{R} \omega_1 \bar{u}}$$

$$\Rightarrow gR\lambda^2 + g \left[ 2\omega_1 \bar{u} + \frac{1}{\omega_1 \bar{u}} - \cancel{g\omega_1 \bar{u}} \right] = 0$$

$$\Rightarrow \boxed{\lambda^2 = - \frac{g}{R} \left[ \omega_1 \bar{u} + \frac{1}{\omega_1 \bar{u}} \right]}$$

Due punti armili di polarizzazione

$$\omega_1 = \sqrt{\frac{g}{R} \omega_1 \bar{u}} \quad \omega_2 = \sqrt{\frac{g}{R} \left( \omega_1 \bar{u} + \frac{1}{\omega_1 \bar{u}} \right)}$$

$$\omega \quad \lambda_{1,2} = \pm i \omega_1 \quad \lambda_{2,1} = \pm i \omega_2$$

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